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THERMODYNAMIC PROPERTIES OF CESIUM UP TO 1500° K)

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THERMODYNAMIC PROPERTIES OF CESIUM UP TO $1500^{\circ}~\mathrm{K}$

By Sheldon Heimel

Lewis Research Center Cleveland, Ohio

NATIONAL AERONAUTICS AND SPACE ADMINISTRATION

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THERMODYNAMIC PROPERTIES OF CESIUM UP TO 1500° K

by Sheldon Heimel

Lewis Research Center

SUMMARY

Consistent tables of thermodynamic properties of elemental cesium were compiled for both the pure species and the equilibrium vapor on the saturation line using selected values of 10 500 calories per mole for the heat of dissociation of dimer and -18 920 calories per mole for the heat of condensation of monomer. The equilibrium vapor properties and properties of condensed cesium (Cs) are given to 1500° K, while properties of gaseous monomer and dimer are given to 2500° K.

The table of mixture properties on the saturation line gives the following thermodynamic functions: temperature T, total vapor pressure P, weight fraction of dimer in vapor \mathbf{x}_2 , molecular weight \mathbf{M}_g , specific volume v, density ρ , enthalpy of condensed phase \mathbf{h}_c , enthalpy of vaporization $\Delta\mathbf{h}_{vap}$, entropy of vapor \mathbf{h}_g , entropy of condensed phase \mathbf{s}_c , entropy of vaporization $\Delta\mathbf{s}_{vap}$, entropy of vapor \mathbf{s}_g , and heat capacities (equilibrium and frozen) $(\mathbf{c}_p)_{eq}$ and $(\mathbf{c}_p)_{fr}$. Assigned enthalpies are relative to enthalpy of the crystal at 298.15° K, which is taken to be zero. The tables of properties for the pure species give the following thermodynamic functions for the standard state: heat capacity at constant pressure \mathbf{C}_p^0 , sensible enthalpy $\mathbf{H}_T^0 - \mathbf{H}_0^0$, entropy \mathbf{S}_T^0 , sensible free energy $\mathbf{F}_T^0 - \mathbf{H}_0^0$, the sum of sensible enthalpy at \mathbf{T}^0 K and chemical energy at \mathbf{H}_T^0 , and values of enthalpy changes and logarithms of the equilibrium constants. The latter two functions are given for the reactions of formation of $\mathbf{C}_{\mathbf{S}}(c)$, $\mathbf{C}_{\mathbf{S}}(c)$, and $\mathbf{C}_{\mathbf{S}}(c)$ from the assigned reference state $\mathbf{C}_{\mathbf{S}}(c)$ (($\Delta\mathbf{H}_T^0$) and $\mathbf{I}_{\mathbf{S}}(c)$) and from $\mathbf{C}_{\mathbf{S}}(c)$

Thermodynamic functions for the gases were generated from atomic and molecular data, whereas the functions of the condensed phase were based on selected experimental data that were smoothed and made self-consistent.

TABLE I. - THERMODYNAMIC PROPERTIES OF CESIUM (CRYSTAL) UP TO 100° K

	H _T - H _O ,	$\mathbf{s_{T}^{o}},$	$-(\mathbf{F}_{\mathbf{T}}^{\mathbf{O}} - \mathbf{H}_{\mathbf{O}}^{\mathbf{O}}),$
	- /		(\- 'T'()' ;
	al/mole	cal/(mole)(^o K)	
5 0.746	1. 2	0.3342	0.5
10 2.379	9.0	1.3539	4.5
15 3.719	23.9	2.5835	14.9
20 4.680	44.5	3.7993	31.5
25 5. 110	69.1	4.8935	53.3
30 5.385	95.3	5.8513	80.2
35 5.558	122.7	6.6955	111.6
40 5.668	150.8	7. 4454	147.0
45 5.741	179.4	8. 1174	185.9
50 5.795	208.2	8.7252	228.1
55 5.842	237.3	9. 2798	273.1
60 5.888	266.6	9. 7901	320.8
65 5.933	296. 2	10. 2631	370.9
70 5.977	325.9	10. 7044	423.4
75 6.016	355.9	11. 1182	477.9
80 6.049	386.1	11.5076	534.5
85 6.075	416.4	11.8751	593.0
90 6.097	446.8	12. 2230	653.2
95 6.122	477.4	12.5533	715.2
100 6.160	508.1	12.8682	778.7

tables III (Cs₁) and IV (Cs₂), data are tabulated at 100° intervals from 0° to 2500° K for the ideal gases. Table V (p. 8) contains the molecular constants used to calculate thermodynamic functions for dimeric cesium.

Table VI (p. 8) contains recommended boiling points of liquid cesium calculated from the final set of thermodynamic properties. These boiling points correspond to pressures from 10⁻⁸ to 10 atmospheres at every power of 10. For comparison, boiling points obtained from equation (33) are also given.

Table VII (p. 9) gives the following thermodynamic properties of the equilibrium gaseous mixture on the saturation line: T, P, p_1 , p_2 , x_2 , M_g , v, ρ , h_c , Δh_{vap} , h_g , s_c , Δs_{vap} , s_g , $(c_p)_{eq}$, and $(c_p)_{fr}$. The values of P, p_1 , p_2 , v, and ρ are given in floating-point notation, where the decimal number is to be multiplied by

10 raised to the power of the sign and the two digits following the letter E; for example, 0.100000E-03 is $0.100000\times10^{-3} = 0.0001$.

Assigned Reference State

The assigned reference state is Cs (crystal, liquid). The crystal is the reference state up to the melting point, and the liquid phase is the reference state above the melting point.

Assigned Enthalpy Values HP

For some applications (ref. 1), it is convenient to combine sensible enthalpy ${\tt H}_T^{\tt O}$ - ${\tt H}_O^{\tt O}$ and chemical energy ${\tt H}_O^{\tt O}$ into one numerical value, ${\tt H}_T^{\tt O}$. The arbitrary base

TABLE II. - THERMODYNAMIC PROPERTIES OF CESTUM (CRYSTAL, LIQUD)

ť, o _M	°Ĉ	H ⁰ - H ⁰ ,	s, T,	$-(F_{\rm T}^{\rm O}-{\rm H_O^{\rm O}}),$	H.T.	Formation from assigned reference state	om assigned e state	Formation from gaseous atoms	on from atoms
	(mole)(°K)	car/ more	(mole)(^O K)	car/ more	car/ more	$(\Delta H_{T}^{0})_{f}$	log ₁₀ K _f	ΔH ^C ,	log ₁₀ K
					,	cal/mole		cal/mole	
0		0	O	0	-1307.2	c	-	-13920.0	1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1
100	9.160	508.1	12.8682	778.7	-1599.1	0	C		36.1543
200	6.560	1144.1	17.2607	2308.0	-063.1	c	2	-18769-5	15.5485
298.15	6.953	1807.2	19.9532	4141.8	0.0	0	t a	-13594.7	7.87.87
300	095*9	1820.1	19.9962	4178.8	12.9	c ·	÷. 1	-18590-3	75.7° 5
3c1.8	6.567	1832.6	20.0379	4514.8	25.4	Ö	C	-18586.8	4.00.0
8361.8	7.455	2352.7	21.7612	4214.8	545.5	9	c .	-14066.7	\$ 0.00 m
400	7.455	3084.8	23.8613	2.6549	1277.6	0	c	-17622.4	5-4664
200	7.455	3450.3	25.5248	8932.1	2023.1	0		-17573.8	\$.5315
007	347 6	4474.8	26.8840	11554-6	2708.6	ာ	c	-17325.1	2.2504
000	4.455	5321.3	28-0332	14301.9	3514.1	c	6	-17076-4	1.3645
000	7.455	6066.8	29.0287	17156.1	4759.6	C	_	-16827.7	7.201.0
000	7.455	6812.3	29.9067	8.60105	5005	c	0	-16575.0	0.1956
b937.71	7.455	7093.4	30.2127	21237.4	5286.2	0	¢.	-16485.2	0.0341
0001	7.455	1557.8	30.6922	23134.4	5750.6	0	c	-16330.3	-0.2341
1100	557 1	8 403.3	31.4027	7.67.49.7	6496.1	c	c	-15081-	-0.5261
1200	7.455	9048-8	32,0514	29412.9	7241.6	0	¢;	-15833.1	-0.7904
1300	7.455	9794.3	32.6481	32648.3	1987.1	0	c	-15584.6	-1.0105
1400	7.455	10539.8	33.2006	35941.0	8732.6	r	ς,	-15336.4	-1.1962
1500	7.455	11285.3	33.7150	39287.1	9478.1	0)	-15088.7	-1.3545

^aMelting point. ^bNormal boiling point to equilibrium mixture.

TABLE III. - THERMODYNAMIC PROPERTIES OF Cs1 (GAS)

T,	ů,	ਸੰ, - ਸੰ,	SO Tr,	-(Fr - HO),	" Н	Formation fr	Formation from assigned	Formation from	on from
<mark>м</mark>	cal	cal/mole	cal	cal/mole	cal/mole	referen	reference state	gaseous atoms	atoms
	(mole)(OK)		(mole)(^O K)			'J(TH∆)	log ₁₀ K _f	$\Delta H_{\mathrm{T}}^{\mathrm{O}},$	log ₁₀ K
						cal/mole		cal/mole	
ď		٥		U	8-21171	0*02681		c	
100	4.9681	496° K	36.5168	3154-9	17609-6	18908.7	-36.1548	. c	c
200	4.5681	993.6	39.9605	6988.5	18105.4	18769.5	-15.5486	c	c
298.15		1481.3	41.9442	11024.4	18594.0	18594.0	-4.9232	c	c
300		1490.4	41.9749	11102.0	18603.2	18590.3	-3.7392	c	r
301.8	4.9681	1489.4	45.0046	11177.6	18612.2	18066.7	-8.6584	C:	c·
8400	4.5681	1987.3	43.4042	15374.4	19100.0	17822.4	-5.4664	c	0
200	4.9681	2484.1	44.5128	19772.3	19596.8	17573.8	-3.5315	0	c
900	4.9681	2580.9	45.4186	24270.2	20093.7	17325.1	-2.2598	c	c
200	4.5682	3477.7	46.1844	28851.4	208507	17076.4	-1.3645	c	0
909	4.9682	3974.5	46.8478	33503.7	21087.3	16827.7	-0.7027	c	C
006	4.9682	4471.3	47.4330	38218.3	21584.1	16579.0	-0.1956	C	Ç
937.71	4.5682	4656.7	47.6369	4001004	217715	16485.2	-0.0341	c,	c
1000	4.9683	4908.2	47.9564	42988.3	22080.9	16330.3	0.2041	0	c ·
1100	4.9686	5465.0	48.4300	47808.0	22577.8	16081.7	0.5261	c	c
1500	4.9654	5961.9	48.8623	52672.9	23074.7	15833.1	0.7904	c	c ·
1300	4.9712	6458.9	49.2601	57579.3	23571.7	15584.6	1.0105	c:	c
1400	6715.4	6950.2	45.6287	62524.0	24069.0	15336.4	1.1962	C	c
1500	4.9814	7454.0	49.9721	67504.2	24566.7	15088.7	1.3545	c	c.
1600	5155*5	7952.6	50.2939	72517.7	25065.4			c	0
1700	5.0079	8452.6	50.5970	77562.3	25565+3			•	C
1800	5.0307	8954.4	50.8839	82636.5	26067.2			0	C:
0061	5.0618	0.6545	51.1566	87738.6	26571.7			c	0
2000	5.1025	1.1965	51.4173	92867.4	27079.9			c	D
2100	5.1541	10479.8	51.6674	98021.8	27592.6			c	c
2200	5.2175	10998.3	51.9086	103200.6	28111.1			C.	Ç.
2300	5.2536	11523.8	52-1422	108403.2	28636.5			c	c
2400	5.3832	12057.5	52.3693	113628.8	29170.3			c. c	c
2500	2.4867	12600.9	52.5911	118876.9	29/13.6			0	
,				•		•			

^aThe reference element is crystal cestum up to 301,8° K and liquid cestum above 301,8° K.

TABLE IV. - THERMODYNAMIC PROPERTIES OF Cs2 (GAS)

0 ¹¹ 0 ¹¹	-	6		(00 02)	moral noticemental on (on oa)	Tormstion fr	Cormetion from sesioned	Formation from	n from
Cp, Hr - Ho, Sr, -	Î E		•	-(FT - HO),	HT.	rormation from assi reference state	ce state	gaseous atoms	atoms
(^o K) (mole)(^o K)	(mole)(OK)					,TH∆)	$\log_{10} \mathrm{K_f}$	$\Delta \mathrm{H}_{\mathrm{T}}^{\mathrm{O}}$	$\log_{10} \mathrm{K}$
						cal/mole		cal/mole	
				C	23725.6	27340.1		-10500.0	
~	_	57,9961		4956*	24568.8	27107.1	-52,3208	-10650.3	19.9889
1743.6	_	64.2330		11103.0	25469.2	26795.5	-22.7862	-10743.6	8,3110
2635.4	_	67.8598		17597.0	26361.0	26361.0	-13.2133	-10827.1	4.4332
_	_	67.9162		17722.6	26377.9	26352.1	-13.0941	-10828.6	4.3847
2668.7	_	01.9707		17844.9	26394.3	25303.3	-12.9797	-10830.0	4.4372
.1893 3567.8		70.5456		24652.0	27293.4	24738-2	-8.5271	-10406-	10040
9.2572 4490.1 72.6074		72.6074		31813.6	28215.7	24169.6	-5.8528	0.87801-	7012-1
9,3259 5419,3 74,3012		74.3012		39161.5	59144.9	23607.7	-4.1114	-11042.4	0.4083
6355.4		75.7441		46665.5	30081.0	23052.8	-2.8967	-11100.0	-0.1578
7298.5		77.0034		54304.2	31024.1	22504.9	-2.0073	-11150.5	-0.6319
8248.9		18.1227		62061.5	31974.5	21964.3	-1.3321	-11193.7	-0.9410
		78.5148		65015.0	32334.8	21762.3	-1.1186	-1120H-1	+040*1-
9.6146 9206.6 79.1317		79.1317		69925.1	32937.2	21431.0	-0,8051	-11229.6	-1.2132
5.6965 10171.8 80.0516	20	80.0516		17884.9	33897.4	20905.3	-0.3843	-11258.1	-1.4366
9.7679 11144.7 80.8981	_	80.8981		85932.9	34870.3	20387.2	-0.0424	-11279.0	-1.6232
5.8468 12125.5 81.6830	_	81.6830		34062.5	35851.1	6.91861	0.2397	-11292.3	-1.7313
9.9273 13114.2 82.4157	_	82.4157		102267.8	36839.8	19374.6	0.4754	-11298.2	-1.9169
10.0053 14111.0 83.1034		83.1034		110544.1	37836.6	18840.4	0.6745	-11296.9	-2.0345
10.0528 15116.1 83.7520		83.7520		118887.2	38841.7			-11289.1	-2.1373
16129.6		84.3664		127293.4	39855.2			-11275.4	-2.2280
17151.7		84.9506		135759.4	40811.3			-11257.1	-2.3085
18182.6	_	85.5080		144282.6	41908.2			-11235.3	-2.3803
10.4425 19222.3 86.0413		86.0413		152860.2	42947.9			-11211-8	-2.4449
10.5337 20271.1 86.5530		86.5530		161490.1	43996.7			-11188.5	-2.5032
21329.1	_	87.0451		170170.2	45054.7			-11167.5	-2.5560
22396.5		87.5196	-11	178898.6	46122.1			-11151.0	-2.6942
23473.3		87.9779	_	187673.6	47198.9			-11141.6	-2.6483
10.9143 24559.9 88.4214	•	88.4214	_	196493.6	48285.5			-11141.8	-2.6889

 $^{
m a}$ The reference element is crystal cesium up to 301.8 $^{
m o}$ K and liquid cesium above 301.8 $^{
m o}$ K.

TABLE V. - MOLECULAR
CONSTANTS FOR DI-

ATOMIC CESIUM

Molecular weight	265.82
Symmetry number	2
Electronic state ^a	1Σ
Statistical weight	1
$\omega_{\rm e}$, cm ^{-1a}	41. 990
$\omega \times .cm^{-1}$	0.08005
$\omega_{\rm o}$ y, cm ⁻¹	-0.0001643
B ₋ . cm ⁻¹⁶	0.01272
$\alpha = \text{cm}^{-1}$	0.000035
D _e , cm ^{-1b}	0.467×10 ⁻⁸

^aConstant from ref. 10.

TABLE VI. - BOILING POINTS OF LIQUID CESIUM FROM 10^{-8} TO 10 ATMOSPHERES

log ₁₀ P _{atm}	Boilin	g points, ^O K
	Recommended	From equation (33) ^a
-8	317.8	317.8
-7	345.8	345.8
-6	379.2	379. 4
-5	420.1	420.3
-4	471. 2	471.4
-3	537.0	537. 2
-2	625.0	625.0
-1	749.1	748.7
0	937. 7	937.3
1	1264.0	1264.5

 $a_{\log_{10}} P_{atm} = \frac{-4053.30}{T} + 7.04453 - 0.915282 \log_{10} T.$

for assigning values to the enthalpy of Cs(c) was a value of zero at 298.15° K. Since table II gives $(H_{298.15}^{O} - H_{O}^{O})_{Cs(c)} = 1807.2$ calories per mole, then $(H_{O}^{O})_{Cs(c)} = -1807.2$ calories per mole. From table II, $(\Delta H_{O}^{O})_{c} = (H_{O}^{O})_{Cs(c)} - (H_{O}^{O})_{Cs_{1}(g)} = -18920.0$ calories per mole. Therefore, $(H_{O}^{O})_{Cs_{1}(g)} = -1807.2 - (-18920.0) = 1712.8$ calories per mole. From table IV, $(\Delta H_{O}^{O})_{2} = -D_{O}^{O} = (H_{O}^{O})_{Cs_{2}(g)} - 2(H_{O}^{O})_{Cs_{1}(g)} = -10500$ calories per mole. Therefore,

$$(H_O^0)_{Cs_2(g)} = -10\ 500.0 + 2(17\ 112.8) = 23\ 725.6$$
 cal/mole

Heats of Formation

Two sets of values for heats of formation are given in tables II to IV. The first set (col. 7) is for the formation of the given species from Cs(c), $(\Delta H_T^O)_f$; the second set (col. 9) is for the formation from $Cs_1(g)$, (ΔH_T^O) . For $Cs_2(g)$ at 298.15° K, for example,

$$(\Delta H_{298.15}^{O})_{f, 2} = (H_{298.15}^{O})_{Cs_{2}(g)} - 2(H_{298.15}^{O})_{Cs(c)}$$
(1)

bConstant from ref. 11.

TABLE VII. - THERMODYNAMIC PROPERTIES OF CESIUM ON SATURATION LINE

											_																														
ρ,	g/1000 cc	0.1134165-34	0. 816142F-08	0.9843095-08	0.1178365-07	0-117837F-37	0.138500F-04	9-9617175-03	0-1525755-01	0.105814F-ng	0.4410725-00	0.1313715 01	0.185815E 01			0.107562E 02			0.3449198 02	(c _p),	.17	cal/(g)(^O K)	0.0373794	0.0373794	0.0373792	0.0373792	0.0373791	0.0373791	0.0373752	0.0373554	0.0373147	0.0372601	0.0372054	0.0371621	0.0371503	0.0371375	9.0371339	0.0371519	0.0371916	0.0372548	0.0373430
٧,	g/oo	.881712E 38				0.848625E 11		0.103981E 07				.761201E 03						•399906E 02	.289923E 02	(c _p) ,	ba .	$\operatorname{cal}/(\operatorname{g})(^{\mathrm{O}}\mathrm{K})$	0.0373794	0.0373801	0.0375824	0.0375969	0.0376118	0.0376118	0.0398111	0.0458544	0.0541423	0.0614825	0.0660140	0.0576452	0.0676643	0.0671774	0.0654981	0.0632733	0.0609153	0.0586486	0.0565811
Mg,	g/mole	132.9100 0.8		0	_	132.9163 0.8		_	_	0	0	0	142.9706 0.5	_	_	_	0	0	55.6196 0.2	S gg		$cal/(g)({}^{\circ}K)$	1,5194868	0.8359656	0.6193402	0.6166767	0.6141195	0.6141193	0.5145876	0.4557130	0.4176329	0.3911611	0.3719515	0.3576685	0.3532508	5.3468938	0.3386943	0.3324155	0.3275886	0.3238748	6.3210207
x2, I	ture	0 000000 132				0.000095 132											_	_	.291861 155	Δs _{vap} ,		cal/(g)(^J K)	1.4226679	0.7060981	0.4692146	0.4662276	0.4633566	0.4503905	0.3350579	0.2636666	0.2153599	0.1802417	0.1535420	0.1326523	0.1259323	0.1159676	0.1024224	0.0912628	0.0319464	0.0740757	0.0673521
×	n g/g															00	0 10	01 0	0 0	S _C ,		cal/(g)(~K)	.0968139	0.1298676	0.1501256	0.1504492	0.1507629	0.1637288	0.1795298	0.1920465	0.2022730		0.2184094		0.2273185	2309262	0.2362719	0.2411527	0.2456422	1661647	.2536686
Partial pressure, atm	p2	0 1636065-2	0.611635E-13	0.804806E-13	0.104799E-12	0.104792E-12	0.296875E-08	0.140152E-05	0.772344E-04	0.126485E-02	0.979553E-02		0.757002E-01	0.155720E-00	0.409726E-00	0.899025E		0	0.466151E	hg,		cal/g ca	.4925 0	136.2305 0.	139.8963 0.		_	_	_	_	_	_			_				0	_	172.3410 0.
Partial pr	p ₁	0.700238E-36	0.150223E-08	0.1823 COE-08	0.219547E-C8	0.219550E-08	0.341672E-05	0.294083E-03	0.549759E-02	0.432016E-01	0.198266E-00		0.924375E 00						0.226204E 02	Δh _{vap} ,		cal/g c	142.2668 132	141.2196 136		_	_	,d				-	, h	-	-	_	_	_			101.0281 172
ъ,	atm	.700238E-36 0	, @	80				3		_		0	01	01	0.1	01	3630E 02	5618E 02		h _c , Δh	_	að	9.7743 142.	1686		_		_			_		32.0488 122.	17	2 :	91	- 1	_		65.7039 103	
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)	100	298	300	a301.	a301.	400	500	9009	700	800	006	D937	1000	1100	1200	1300	1400	1500				10	20	29	300	301	4301,	400	200	900	200	900	b 200	786~	0001	1100	1200	1300	0047	150

 $^{^{\}mathrm{a}}_{\mathrm{Melting\ point.}}$ bornmal boiling point to equilibrium mixture.

$$(\Delta H_{298.15}^{O})_{2} = (H_{298.15}^{O})_{Cs_{2}(g)} - 2(H_{298.15}^{O})_{Cs_{1}(g)}$$
 (2)

Equilibrium Constants

Two sets of logarithms of the equilibrium constants for the two formation reactions discussed previously are also listed in tables II to IV. The equilibrium constant K_f for formation from the assigned reference element (col. 8) is obtained from the standard free energy change $(\Delta F_T^0)_f$ by means of the equation

$$\log_{10} K_{f} = -\frac{(\Delta F_{T}^{O})_{f}}{2.3025851 \text{ RT}}$$
(3)

The equilibrium constant K for formation from the atomic gas is obtained from a similar equation

$$\log_{10} K = -\frac{\Delta F_{\rm T}^{\rm O}}{2.3025851 \text{ RT}} \tag{4}$$

Atomic weights, the universal gas constant and the constants used in the evaluation of the entropy constant were the same as those used in reference 1.

SELECTION OF INITIAL DATA

Crystal

Heat capacities of the crystal (table I, p. 4) were derived by smoothing the experimental data from the sources given in table VIII.

The following procedure was used to obtain smooth data from 0° to 100° K. A least-squares fit of the data of reference 2 was used in the temperature range from 0° to 4° K. Data from 4° to 20° K were obtained from a curve drawn through the data of reference 2 at 4° K and through the data of references 3 and 4 at 20° K. Data from 20° to 100° K were obtained from a least-squares fit of the data of reference 4 in that temperature range.

Between 100° and 210° K, the data of reference 4 are approximately linear in tem-

TABLE VIII. - REFERENCES FOR HEAT

CAPACITIES OF CRYSTAL

Temperature range,	Source
0.3 to 4 8 to 12	Ref. 2 Private communication from Professor D. C. McCollum of University of California in Riverside, California and ref. 3
20 to 210	Ref. 4

perature. The heat capacities at 100° and 210° K yield the equation

$$C_p^0 = 5.76 + 0.004 \text{ T}$$
 (5)

Above 210° K, thermal effects associated with the cesium-oxygen system cause anomalous increases in heat capacity and therefore these data were not used. Equation (5) was used to extrapolate heat capacities up to the melting point. The melting point of 301.8° K and the

heat of fusion of 520.1 calories per mole were taken from reference 5.

Reference 4 gives a melting point of 300.5° K and the total heat absorption between the crystal at 262.5° K and the liquid at 300.5° K. If this heat is compared with the sensible enthalpy change of the crystal between these two temperatures, as calculated from equation (5), the implied heat of melting is 514 calories per mole. This value is in reasonable agreement with the value adopted.

The enthalpy from the linear equation for C_p^0 , when combined with the heat of fusion, gave

$$(H_{301.8}^{O})_{Cs(liquid)} - (H_{298.15}^{O})_{Cs(crystal)} = 545.5$$
 cal/mole (6)

Liquid

A number of empirical equations have been used to satisfactorily fit enthalpy data as a function of temperature for various substances; however, heat capacity data derived from differentiation of the enthalpy equations are often very unreliable. For any particular species, various empirical enthalpy equations that might be chosen yield considerably different sets of heat capacities.

Two sets of experimental enthalpy data were considered: that of references 6 and 7. The curve of reference 6 exhibits an anomalous bump from 351° to 620° K, followed by a straight line to 1238° K. It is recommended in reference 6 that just the straight-line portion be used, which has a standard deviation $\sigma = 1.5$ percent. For the cubic equation (from 340° to 1176° K) of reference 7, $\sigma = 1.1$ percent.

The implied $C_p^0 = 7.25$ calories per mole per 0K of reference 6 differs greatly from the heat capacities derived from the enthalpy curve of reference 7 and listed in table IX.

According to reference 7, the curves of heat capacity against temperature of liquid

TABLE IX. - HEAT CAPACITIES

DERIVED FROM ENTHALPY

CURVE OF REFERENCE 7

Absolute	Heat capacity at
temperature,	constant pressure for
T,	standard state,
°К	C _p ,
	cal/(mole)(^O K)
301.8	10. 35
500	7. 28
700	6.23
1000	8.53
1200	12.64
1500	22.67

cesium and other alkali metals have a parabolic shape; that is, heat capacity initially decreases with increasing temperature, reaches a minimum, and then increases continuously. However, a parabolic shape for heat capacity of these species is only the result of selecting a cubic to represent experimental enthalpy data.

It was thought advisable to use a straight line for enthalpy because (1) the straight line, in the standard deviation sense, represents cesium enthalpies about as well as the cubic does and (2) the cubic has the disadvantage of implying very large heat capacities for even a few hundred degrees extrapolation.

The actual straight line used in this report was based on the straight-line data of reference 6 but was

constrained to go through the previously calculated relative heat content of 545.5 calories per mole at the melting point (eq. (6)). This gives the following equation:

$$(H_T^0)_{Cs(liquid)} - (H_{298.15}^0)_{Cs(crystal)} = -1704.5 + 7.455 T$$
 (7)

Equation (7) implies a constant value of $C_p^0 = 7.455$ calories per mole per 0K .

Had the data of reference 7 been fitted to a straight line and similarly constrained to go through the heat content of 545.5 calories per mole, a constant value of $C_p^0 = 7.595$ calories per mole per 0K would have resulted. The enthalpies would be, on the average, only about 1 percent higher than those derived from the data of reference 6.

The enthalpy data of reference 8 are about 6 to 7 percent lower than the data of both references 6 and 7 over the range from about 600° to 1300° K; therefore, they were not used.

Monomer (Gas)

The partition function used to compute the thermodynamic properties of Cs₁ was truncated by the temperature-dependent cutoff technique used in reference 1. The results of these computations are exactly the same up to 2200° K as would be obtained by simply using all levels given in reference 9 with no cutoff. Above 2200° K enthalpy and entropy differ only in the last place computed and heat capacity in the last or second last place.

Dimer (Gas)

The method used to compute the thermodynamic properties of Cs_2 is the method for diatomic molecules described in reference 1. The constants for the dimer are shown in table V(p. 8). They were taken from references 10 and 11.

Vapor Pressures

During the period 1913 to 1937, the investigations reported in references 12 to 18 obtained vapor pressure data for cesium in the temperature range from 238° to 670° K. This early work was evaluated in reference 19 where it was concluded that, except for the results of references 17 and 18, the early work is not too reliable.

The vapor pressures in the 455° to 589° K range (ref. 17) were used in reference 19 to generate a vapor pressure equation. Inasmuch as the data of reference 17 are only relative pressure data (obtained by the magneto-optical method) and are converted to absolute pressure data by relating them to the unreliable data of reference 13, it seems prudent to disregard the data of reference 17.

Recently, the first experimental measurements in the temperature range that included the boiling point $(742^{\rm O}\ {\rm to}\ 1199^{\rm O}\ {\rm K})$ were reported in reference 20. These data check reasonably well with the later data of reference 6 $(727^{\rm O}\ {\rm to}\ 1334^{\rm O}\ {\rm K})$.

In view of the previous observations, the data of references 18 (vapor pressures in the liquid range), 20, and 6 were combined and fitted by the least-squares technique. All the data were given the same weight. The results of several empirical equations were compared. The following equation, which is in the form of the Kirchhoff equation, was accepted because it gave as satisfactory results as the other forms considered:

$$\log_{10} P = -\frac{3920.38}{T} + 5.71342 - 0.519781 \log_{10} T$$
 (8)

If only vapor pressures were considered in this report, then the smoothed experimental vapor pressures obtained from equation (8) would be accepted. As pointed out in the INTRODUCTION, however, these smoothed data, when combined with other thermodynamic data, lead to inconsistencies in the heat of condensation of the monomer at $0^{\rm O}$ K ($\Delta {\rm H}_{\rm O}^{\rm O}$). Therefore, the data from equation (8) are used only for preliminary analysis, whereas the final recommended vapor pressures are calculated as explained in detail in the section Pressures and Weight Fractions of Gaseous Species.

Heat of Dissociation at 00 K

An upper bound (4020 cm $^{-1}$) for the heat of dissociation at $0^{\rm O}$ K(${\rm D_O^{\rm O}}$) is estimated in reference 21 by analytical extrapolation of the ground state vibrational intervals. A lower bound (3318 cm $^{-1}$) is estimated in reference 21 by making certain assumptions as to the final states of the excited atoms produced. The average of these two values, which is equivalent to 10 500±1000 calories per mole of dimer, is recommended in this reference.

The estimate of 0.453 eV or 10 450 calories per mole from the correlation of force constants, equilibrium internuclear distances, and heats of dissociation of the alkali metals is obtained in reference 22. A technique for deducing a value of D_O^O where D_O^O was not well known is suggested in reference 23. This technique gave an optimum D_O^O = 10 000 calories per mole which is in fairly close agreement with the spectroscopic value of 10 500 calories per mole.

The calculation of $(\Delta H_O^O)_c$, which is discussed in the next section, depends in part on the value of D_O^O . While no value of D_O^O that was tried gave a constant value of $(\Delta H_O^O)_c$ over the entire temperature range of 301.8° to 1500° K, values of D_O^O from about 10 000 to 11 500 calories per mole gave the most nearly constant value of $(\Delta H_O^O)_c$. Based on this analysis, the conclusion was that the spectroscopic value was approximately correct, and the value $D_O^O = 10$ 500 calories per mole was selected.

Heat of Condensation of Monomer at 0^{0} K

The values of $(\Delta H_O^O)_c$ in calories per mole are given in a number of publications as follows: -19 048 (ref. 24), -18 790 (ref. 19), -19 050 (ref. 25), and -19 035 (ref. 26). The aforementioned values of $(\Delta H_O^O)_c$ are based on the pre-1962 vapor pressures. Therefore, this report will recommend a ''best'' (ΔH_O^O) of -18 920 on the basis of the newer vapor pressures, as discussed in the following section.

THERMODYNAMIC ANALYSIS

As discussed in the INTRODUCTION, the procedure for the selection of a value for $(\Delta H_O^0)_c$ is to find that $(\Delta H_O^0)_c$ for which the difference between the calculated and smoothed experimental vapor pressures is a minimum. As a first step, it is necessary to compute $(\Delta H_O^0)_c$ for each temperature in the range of interest.

Derivation of $(\Delta H_0^0)_C$

The zero-degree heat of condensation of monomer was computed by means of the relation

$$(\Delta H_O^0)_c = (\Delta F_T^0)_c - \Delta (F_T^0 - H_O^0)_c$$
(9)

where

$$(\Delta F_T^0)_c = (F_T^0)_{Cs(c)} - (F_T^0)_{Cs_1(g)} = RT \ln p_1$$
 (10)

and

$$\Delta(F_{T}^{o} - H_{O}^{o})_{c} = (F_{T}^{o} - H_{O}^{o})_{Cs(c)} - (F_{T}^{o} - H_{O}^{o})_{Cs_{1}(g)}$$
(11)

Partial Pressure of Monomer

The partial pressure of monomer $\,p_1^{}$ in atmospheres is obtained from the simultaneous solution of Dalton's Law and the equilibrium-constant equation for the dimerization reaction

$$2 \text{ Cs}_1 = \text{Cs}_2$$

The equations are

$$P = p_1 + p_2 \tag{12}$$

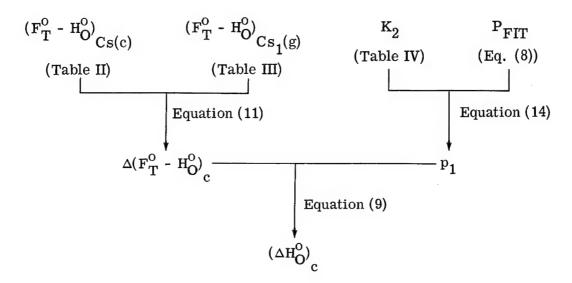
and

$$K_2 = \frac{p_2}{p_1^2} \tag{13}$$

where P is the total pressure, p_2 is the partial pressure of dimer, and K_2 is the equilibrium constant of dimerization. Values of K_2 are obtained from table IV (p. 7). The solution of equations (12) and (13) is

$$p_1 = \frac{-1 + \sqrt{4PK_2 + 1}}{2K_2} \tag{14}$$

Therefore, apparent values of $(\Delta H_O^0)_c$ can be computed at every temperature of interest by starting with given values of D_O^0 , values of F_T^0 - H_O^0 for monomer, dimer, and condensed phase, and smoothed experimental vapor pressures P_{FIT} . This procedure can be summarized as follows:



RESULTS OF CALCULATING (ΔH_0^0)

The apparent $(\Delta H_O^0)_c$ values computed by means of this procedure from equation (9) are given in table X. The maximum difference is about 1 percent.

If some single arbitrary value of $(\Delta H_O^0)_c$ is selected, then for each of the previous temperatures there will be a difference between the smoothed experimental vapor pressure P_{FIT} and the vapor pressure P_{CALC} calculated with this arbitrary $(\Delta H_O^0)_c$ value.

For this analysis, an error function may be defined as $|P_{FIT} - P_{CALC}|/P_{FIT}$ averaged over the previous number of temperatures. The procedure for selecting a value of $(\Delta H_O^0)_c$ is to find that value of $(\Delta H_O^0)_c$ which gives the minimum error function.

The results of calculating the error function for various values of $(\Delta H_O^0)_c$ are given in table XI. The optimum $(\Delta H_O^0)_c$ was taken to be -18 920 calories per mole to the nearest 5 calories.

The fact that there is still about 4 percent error between the smoothed experimental vapor pressure and vapor pressure calculated with the optimum $(\Delta H_O^0)_c$, indicates that

Table X. - Apparent values of -($\Delta H_{\mathrm{O}}^{\mathrm{o}}$) $_{\mathrm{c}}$

AS A FUNCTION OF TEMPERATURE

Tempera- ture, T,	-($\Delta H_{O}^{0})_{c}$, cal/mole	Tempera- ture, T, OK	$-(\Delta H_{ m O}^{ m O})_{ m c}$, cal/mole
301.8	18 791.6	1000	18 965.1
400	18 872.5	1100	18 954.3
500	18 921.7	1200	18 937.0
600	18 949.3	1300	18 913.1
700	18 963.6	1400	18 881.9
800	18 969.7	1500	18 843.4
900	18 970.0		

TABLE XI. - ERROR FUNCTION

FOR VARIOUS VALUES

OF
$$-(\Delta H_O^O)_C$$

$-(\Delta H_{\mathrm{O}}^{0})_{\mathrm{c}}^{},$ cal/mole	Error function, ΔP /P
18 850	0. 0538
18 900	. 0397
18 910	. 0381
18 920	. 0370
18 925	.0370
18 930	. 0372
18 950	. 0389
18 970	.0458

discrepancies still remain which may be in any of the following data:

- (a) Experimental vapor pressures
- (b) D_O
- (c) Molecular constants of dimer (and, consequently, its thermodynamic properties)
- (d) Heat content (and other derived thermodynamic properties) of the condensed phase
- (e) $(\Delta H_O^0)_c$

The 4 percent error may also be due, in part, to the presence of a higher polymer, for example, trimer or tetramer; however, this possibility has not been considered in this report.

CALCULATION OF EQUILIBRIUM MIXTURE PROPERTIES ON SATURATION LINE

As explained in the INTRODUCTION, selection of an optimum value of $(\Delta H_O^o)_c$ makes it possible to compute a set of vapor pressures consistent with the sensible free energies of the pertinent species. These vapor pressures and other equilibrium properties on the saturation line (e.g., enthalpy, specific heat, and entropy) are given in table VII (p. 9). The equations used to obtain these properties are given in the following sections.

Pressures and Weight Fractions of Gaseous Species

By starting with the optimized $(\Delta H_O^0)_c$ and the sensible free energies of monomer and condensed phase, monomer pressure was computed with equations (9) to (11).

Dimer pressure was computed from equation (13). The total pressure P_{CALC} was obtained from the sum of p_1 and p_2 .

The weight fraction of monomer x_1 was obtained from the mole fraction of monomer N_1 ; thus,

$$N_1 = \frac{P_1}{P} \tag{15}$$

$$x_1 = \frac{N_1}{2 - N_1} \tag{16}$$

The weight fraction of the dimer x2 was obtained from the defining equation

$$x_1 + x_2 = 1 (17)$$

Enthalpy of Vaporization and Entropy of Vaporization

The heat of vaporization in calories per gram from the condensed state to the equilibrium mixture on the saturation line was obtained from

$$\Delta h_{\text{vap}} = \frac{2x_1(\Delta H_T^0)_{f, 1} + x_2(\Delta H_T^0)_{f, 2}}{M_2}$$
 (18)

where M_2 is the molecular weight of dimer (265.82 g/mole). The heats of formation are obtained from tables III and IV (pp. 6 and 7).

The entropy of vaporization in calories per gram is obtained from equation (18) by means of

$$\Delta s_{\text{vap}} = \frac{\Delta h_{\text{vap}}}{T} \tag{19}$$

Enthalpy, Specific Heat, and Entropy of Cesium Vapor

The enthalpy $h_{\mathbf{g}}$ in calories per gram of an equilibrium mixture can be defined by

$$h_{g} = \sum_{i=1}^{2} \frac{x_{i}(H_{T}^{O})_{i}}{M_{i}} = \sum_{i=1}^{2} x_{i}h_{i}$$
 (20)

Equation (20) for saturation conditions gives the same results as the usual equation for the enthalpy of a mixture

$$h_g = h_{vap} + h_c$$

where h_c is the enthalpy of the condensed phase in calories per gram (table VII, p. 9). The equilibrium specific heat at constant pressure in calories per gram per ${}^{O}K$ is defined by

$$(c_p)_{eq} = \left(\frac{\partial h}{\partial T}\right)_p$$
 (21)

Using equation (20) results in equation (21) becoming

$$(c_p)_{eq} = \sum_{i=1}^{2} x_i (c_p)_i + \sum_{i=1}^{2} h_i \left(\frac{\partial x_i}{\partial T}\right)_p$$
 (22)

or

$$(c_p)_{eq} = (c_p)_{fr} + (c_p)_r$$
(23)

From equation (17), it follows that

$$\left(\frac{\partial \mathbf{x}_1}{\partial \mathbf{T}}\right)_{\mathbf{p}} = -\left(\frac{\partial \mathbf{x}_2}{\partial \mathbf{T}}\right)_{\mathbf{p}} \tag{24}$$

Therefore,

$$(c_p)_r = \left(\frac{\partial x_2}{\partial T}\right)_p (h_2 - h_1) = \left(\frac{\partial x_2}{\partial T}\right)_p \frac{(\Delta H_T^0)_2}{M_2}$$
 (25)

Values of $(\Delta H_T^0)_2$ are tabulated in table IV (p. 7). Values of $(\partial x_2/\partial T)_p$ may be calculated from the following equation:

$$\left(\frac{\partial x_2}{\partial T}\right)_p = \frac{(\Delta H_T^0)}{RT^2} \frac{x_2(2 - x_2)(1 - x_2)}{2}$$
(26)

Equation (26) may be derived from the following expression, which is equivalent to equation (13):

$$K_2 = \frac{x_2(2 - x_2)}{4P(1 - x_2)^2} \tag{27}$$

and from the van't Hoff isobar, namely,

$$\left(\frac{\partial \ln K_2}{\partial T}\right)_p = \frac{(\Delta H_T^0)_2}{RT^2}$$
 (28)

The entropy of the gas mixture may be obtained from

$$s_g = \Delta s_{vap} + s_c \tag{29}$$

where s_c is the entropy of the condensed phase in calories per gram per ${}^{O}K$ (table VII, p. 9) and Δs_{vap} is obtained from equation (19). Equation (29) gives the identical results as the usual equation for the entropy of a mixture.

Molecular Weight of Mixture

An expression for the molecular weight of the mixture M_g (in g/mole) in terms of the molecular weight and weight fraction of the dimer may be obtained as follows:

$$M_{g} = \frac{1}{n} = \frac{1}{n_{1} + n_{2}} = \frac{1}{\frac{x_{1}}{M_{1}} + \frac{x_{2}}{M_{2}}} = \frac{M_{2}}{2 - x_{2}}$$
(30)

where n, n_1 , and n_2 are the number of moles of mixture, monomer, and dimer per gram of mixture, respectively.

Specific Volume and Density of Mixture

The specific volume $\,v\,$ in cubic centimeters per gram may be obtained from the ideal gas law

$$v = \frac{RT}{PM_g} = \frac{82.05971 \text{ T}}{PM_g}$$
 (31)

The density ρ in grams per 1000 cubic centimeters is, therefore,

$$\rho = \frac{1000}{v} \tag{32}$$

DISCUSSION OF VAPOR PRESSURES

Vapor Pressures Above Liquid

Table VI contains boiling points of liquid cesium at saturation pressures from 10⁻⁸ to 10 atmospheres at every power of 10. These data were calculated in two ways:

- (1) From the recommended thermodynamic data, as explained in the section THERMODYNAMIC ANALYSIS
- (2) By a least-squares equation fitted to the recommended P_{CALC} values with a standard deviation of 0.5 percent

$$\log_{10} P_{\text{atm}} = \frac{-4053.30}{T} + 7.04453 - 0.915282 \log_{10} T$$
 (33)

On the average the two sets of boiling points differ by about 0.2° K. It should be emphasized that the vapor pressure equation (eq. (33)) is recommended rather than equation (8).

Figures 1 and 2 compare the experimental vapor pressures of liquid cesium with a curve drawn through the recommended vapor pressures. The agreement of the two sets is seen to be very good. Figure 1 covers the temperature range from 303° to 1334° K. Figure 2 covers only the range from 729° to 1334° K to facilitate detailed comparison.

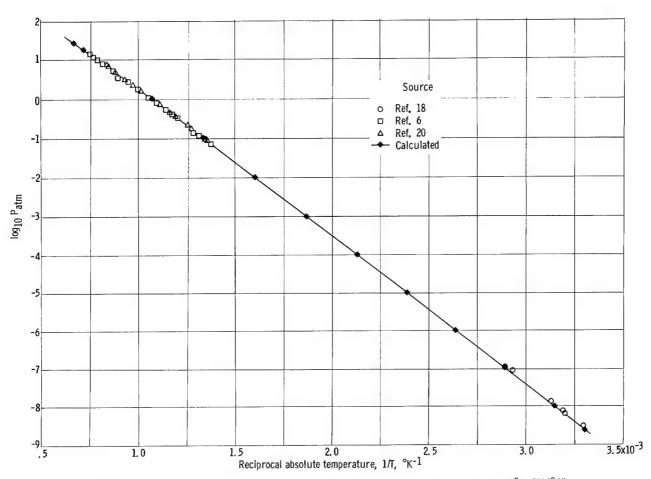


Figure 1. - Comparison of experimental and calculated vapor pressures above liquid cesium (303 $^\circ$ to 1334 $^\circ$ K).

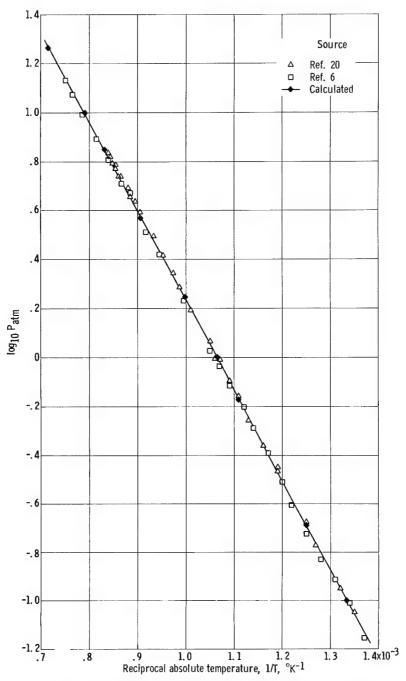


Figure 2. – Comparison of experimental and calculated vapor pressures above liquid cesium (729° to 1334° K).

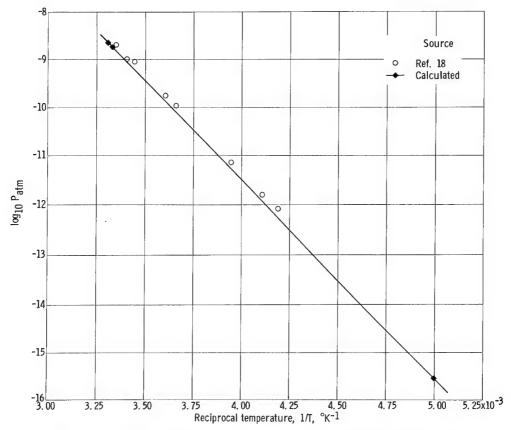


Figure 3. - Comparison of experimental and calculated vapor pressures above crystalline cesium.

Subsequent to the completion of the analysis, a vapor pressure equation in the temperature range 928° to 1558° K was presented in reference 27. On the average, these vapor pressures are about 2 percent lower than those recommended in this report.

Vapor Pressures Above Crystal

Experimental vapor pressures above the crystal are available from reference 18 at temperatures from 238° to 298° K. In order to be consistent with the thermodynamic data adopted for this report, however, vapor pressures from 100° to 301.8° K were calculated in the same manner as were the vapor pressures above the liquid.

On the average, the calculated vapor pressures between $238^{\rm O}$ and $298^{\rm O}$ K were 30 percent lower than the experimental ones (fig. 3). This difference is not significant, however, inasmuch as at the low pressure levels involved, the average difference is only in the order of 10^{-9} atmosphere. Furthermore, the lower calculated vapor pres-

sures are consistent with the assumption of reference 26 that the experimental vapor pressures are too high.

Lewis Research Center,
National Aeronautics and Space Administration,
Cleveland, Ohio, April 26, 1965.

APPENDIX - SYMBOLS

B _e C _p ^o	rotational constant, cm ⁻¹ heat capacity at constant pressure for standard	$(\Delta H_{\mathrm{T}}^{\mathrm{O}})_{\mathrm{f}}$	enthalpy change for formation of substance from assigned reference element, cal/mole
	state, cal/(mole)(^o K)	$H_{\mathbf{O}}^{\mathbf{O}}$	chemical energy at 0° K for standard state, cal/mole
$(c_p)_{eq}$	equilibrium specific heat at constant pressure of mixture, cal/g	h	enthalpy, cal/g
		$^{\Delta h}$ vap	heat of vaporization to 1 gram of equilibrium vapor
$(c_p)_{fr}$	frozen specific heat at con- stant pressure of mixture, cal/g	K	equilibrium constant for re- action of formation from element in atomic gas state
(c _p) _r	reactive specific heat at con- stant pressure of mixture, cal/g	K_{f}	equilibrium constant for reaction of formation from
D_{e}	spectroscopic constant for rotational stretching, cm ⁻¹	M	assigned reference element molecular weight based on
D_{O}^{O}	dissociation energy at 0° K, cal/mole		chemical scale of natural oxygen, g/mole
$\mathbf{F_{T}^{o}}$	Gibbs free energy for stand- ard state, cal/mole	N _i	mole fraction of i th species, moles of i th species per mole of mixture
F_{T}^{O} - H_{O}^{O}	sensible free energy for standard state, cal/mole	n	number of moles per gram of equilibrium mixture
$H_{\mathbf{T}}^{\mathrm{O}}$	sum of sensible enthalpy at T ^O K and chemical energy	P	total vapor pressure, atm
	at 0° K for standard state,	$\mathbf{p_i}$	partial pressure of i th species, atm
$H_{\mathrm{T}}^{\mathrm{O}}$ - $H_{\mathrm{O}}^{\mathrm{O}}$	sensible enthalpy for stand- ard state, cal/mole	R	universal gas constant, 1. 98726 cal/(mole)(^O K) or 82. 05971 (cc)(atm)/
$\Delta H_{\mathrm{T}}^{\mathrm{O}}$	enthalpy change for forma-		(mole)(^o K)
	tion of substance from ele- ment in atomic gas state, cal/mole	$\mathbf{s}_{\mathrm{T}}^{\mathrm{o}}$	entropy for standard state, cal/(mole)(^O K)

s	entropy, cal/(g)(OK)	$^{\omega}\mathrm{_{e}}$	zero-order vibrational fre-
$\Delta s_{ m vap}$	entropy of vaporization to 1 gram of equilibrium vapor		quency for diatomic molecule, cm ⁻¹
T	absolute temperature, ^O K	$\omega_{e}^{x}_{e}, \omega_{e}^{y}_{e}$	anharmonicity constants for diatomic molecule, cm ⁻¹
v	specific volume of vapor, cc/g	Subscripts:	
x_i	weight fraction of i th species, g of i th species per g of	c	condensed phase property
	mixture	f	formation from assigned
$^{lpha}_{ m e}$	vibration-rotation interaction		reference element
	constant for diatomic mole- cule, cm ⁻¹	g	property of equilibrium vapor mixture
ρ	density of vapor, g/1000 cc	1	monomer property
σ	standard deviation	2	dimer property

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